

Mathematical modelling of the loss of tissue compression responsiveness and its role in solid tumour development

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This paper presents a mathematical model of normal and abnormal tissue growth. The modelling focuses on the potential role that stress responsiveness may play in causing proliferative disorders which are at the basis of the development of avascular tumours. In particular, we study how an incorrect sensing of its compression state by a cell population can represent a clonal advantage and can generate hyperplasia and tumour growth with well-known characteristics such as compression of the tissue, structural changes in the extracellular matrix, change in the percentage of cell type (normal or abnormal), extracellular matrix and extracellular liquid. A spatially independent description of the phenomenon is given initially by a system of non-linear ordinary differential equations which is explicitly solved in some cases of biological interest showing a first phase in which some abnormal cells simply replace the normal ones, a second phase in which the hyper-proliferation of the abnormal cells causes a progressive compression within the tissue itself and a third phase in which the tissue reaches a compressed state, which presses on the surrounding environment. A travelling wave analysis is also performed which gives an estimate of the velocity of the growing mass.

Keywords: tissue growth; stress responsiveness; proliferative disorders; avascular tumour growth.

1. Introduction

Over the last few years, there has been a lot of attention paid to modelling avascular tumour growth and tumour invasion of host tissue, as can be seen from the books (Adam & Bellomo, 1996; Preziosi, 2003), the special issues of journals (Bellomo & De Angelis, 2003; Chaplain, 1999, 2002) and review articles (Araujo & McElwain, 2004a) specifically devoted to these topics, and the references therein. Most of the modelling carried out in these papers starts from the observation that tumour cells become insensitive to anti-growth signals and evade apoptosis. Having lost their internal programme-for-cell-death, they acquire a limitless replicative potential, grow rapidly and invade the surrounding tissue (see Araujo & McElwain, 2004a; Bellomo *et al.*, 2003; Chaplain & Anderson, 2003, for recent reviews). Attention is then focussed on how the nutrient (e.g. oxygen) diffusing through the tumour boundary affects the growth of the tumour, with the formation of a necrotic core, of an intermediate layer of quiescent cells

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and an outer rim of proliferating cells. Most mathematical models are characterized by the existence of a limiting radius, as is well-known from experiments (Freyer & Sutherland, 1986; Sutherland & Durand, 1984).

In this paper, we start from a different point of view and examine certain important biomechanical factors. Specifically, we focus on the role of how a cell senses the presence of other cells and its own compression state, how this affects its proliferative ability and how an imperfect perception may cause hyperplasia, which here is considered as an early event of tumorigenesis. Thus, for the purposes of our model, we neglect the effects of external nutrients and other chemical/growth factors and choose to focus solely on biomechanical effects.

As will be discussed in Section 2, which will focus on the description of experimental evidence, several papers have studied the effect of stress on the process of mitosis and apoptosis of tumour cells. In order to deal with growing biological tissues, recently the concepts of multiple natural configurations and of accretive forces have been introduced to model tissues as continua which continuously remodel while undergoing large deformations (Ambrosi & Guana, 2005; Ambrosi & Mollica, 2002, 2003; Di Carlo & Quiligotti, 2002; Fusi *et al.*, 2006; Humphrey & Rajagopal, 2002; Rao *et al.*, 2003). In fact, the simultaneous presence of growth and deformation makes it impossible to use standard continuum mechanics. This approach was used by Ambrosi & Mollica (2004) to model the growth in vitro of a multicellular spheroid in gels of variable stiffness obtaining results in good qualitative agreement with the experiment by Helmlinger *et al.* (1997).

On the other hand, to take into account the different constituents present in a tumour, the use of the theory of mixtures was proposed in Byrne & Preziosi (2004). This approach also enables the stresses within the tumour to be evaluated. The link between the model proposed here and the one in Byrne & Preziosi (2004) consists in the inclusion of the dependence of the cell proliferation rate on the cellular stress. Many other applications of multiphase models have been proposed in Ambrosi & Preziosi (2002), Araujo & McElwain (2004b), Breward *et al.* (2001, 2002, 2003), Byrne *et al.* (2003), Franks & King (2003) and Preziosi & Graziano (2003), some of them also including the effect of the mechanical interactions with the surrounding tissues.

The main aim of this paper is to show how an underestimation of the compression state of the local tissue, and then of the subsequent stress which is exerted on a cell, can generate by itself a clonal advantage on the surrounding cells leading to the replacement and the invasion of the healthy tissue.

In addition to biomechanical effects, the model also takes into account the effect of the production of extracellular matrix (ECM) and of matrix-degrading enzymes (MDEs). In fact, cells produce ECM which is important for cell adhesion, spreading and motility. Therefore, on the one hand the ECM may constitute a barrier to normal cell movement, and on the other hand it provides a substratum cells may use to adhere and move. Most mammalian cell types require at least some elements of the ECM to be present for growth and survival and will indeed migrate up a gradient of bound (i.e. non-diffusible) cell adhesion molecules in culture in vitro (Carter, 1965; Klominek *et al.*, 1993; Lacovara *et al.*, 1984; Lawrence & Steeg, 1996; McCarthy & Furcht, 1984; Quigley *et al.*, 1983).

On the other hand, MDEs are also important at many stages of tumour growth, invasion and metastasis, and the manner in which they interact with inhibitors, growth factors and tumour cells is very complex. However, it is well-known that tumour cells produce MDEs which degrade the ECM.

In summary, we start from the assumption that when cells are in a crowded environment they sense the presence of other cells and their behaviour then crucially depends on how they can withstand the pressure. We then focus on how this can affect both mitosis and production of ECM and MDEs. We start from the following phenomenological observations which are described in greater detail in Section 2.

- Cells replicate if they sense there is sufficient space to accomplish this. If, on the other hand, they sense that there is a sufficient number of cells around them, they can alter their activity and enter a quiescent state ready to re-activate their replication programme if, for instance, a neighbouring cell dies.
- Cells constantly produce ECM and MDEs.
- Cells move preferentially towards regions with lower stress.

Of course, we are well aware that cellular mechano-transduction is not the only cause of the formation of hyperplasia and tumours, and that chemical factors will operate to regulate the reproduction rates. However, in this paper we want to focus specifically on the role of cell contact and compression. In fact, we assume that the only thing that changes a normal cell into an abnormal (transformed) cell is how it senses and responds to the stress exerted on it. If a cell does not sense stress correctly, it continues to replicate even when there are already other cells around it and insufficient space for the daughter cells. This gives rise to hyperplasia and tumours in the local tissue. In this paper then, the stress and therefore its influence on the evolution of the cell population occurs through three contributions: cell replication, the production of ECM and the release of MDEs.

The paper is laid out as follows. Section 2 is devoted to a detailed description of the biological aspects covered by the model developed in Section 3. We then study in Section 4 the properties of the model starting from the simplest version and then including more complicated effects along the way, giving analytic solutions and qualitative behaviours whenever possible. More precisely, Section 4.1 describes how the model is able to describe the generation of confluent cell distributions in vitro in view of the identification of some parameters. Section 4.2 focuses on the re-construction of a normal healthy tissue with MDEs controlling the content of the ECM and allowing a continuous renewal of the ECM. Section 4.3 describes how the occurrence of a genetic change in a cell causing only an inadequate sensing of stress which affects only cell replication can lead to the complete replacement of the healthy tissue and to an overcompression of the tissue. In particular, the normal ECM is completely replaced by the tumour-generated matrix. Section 4.4 describes what happens in the case of a hyper-production or hypo-production of ECM or MDEs by tumour cells. This incorrect balance is, for instance, characteristic of fibrosis (Johnson, 2001). Finally, Section 5 considers the spatially dependent situation showing how the phenomena described from an analytical point of view in Section 4 manifest themselves in a spatial domain where, however, the overcompression of hyperplasia with respect to the healthy tissue causes the progressive invasion of the tissue. A qualitative description of the travelling wave solution is also given. The velocity of invasion increases with the level of incorrect sensing of stress, and in particular it is proportional to the square root of the difference of the volume ratio at which abnormal and normal cells stop duplicating.

2. Phenomenological description

It is well-known that the rate of proliferation of most cells decreases when they come in contact with other cells. This phenomenon is often called contact inhibition of growth (Deleu *et al.*, 1998; Dietrich *et al.*, 1997; Kato *et al.*, 1997; Nelson & Chen, 2003; Polyak *et al.*, 1994a; St Croix *et al.*, 1998). A quantification of this phenomenon is represented in Fig. 1 that reports some experimental results by Tzukatani *et al.* (1997) on human breast epithelial cells grown in vitro over a suitable substratum. It can be seen that after an initial exponential growth, cell density saturates.

Several experimental papers describe how the initiators of the phenomenon are the cadherins, the transmembrane receptors involved in homophilic cell–cell interactions, because of their crucial role

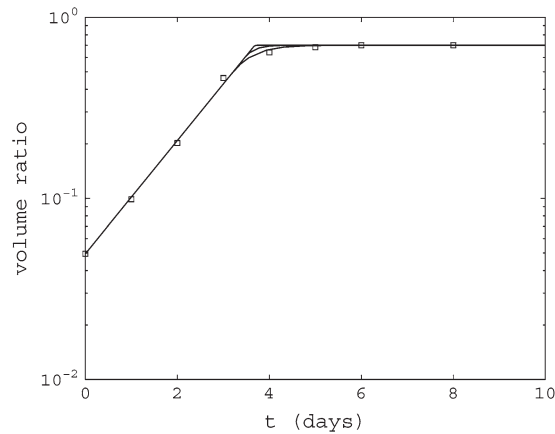


FIG. 1. Growth of human breast epithelial cells. The squares refer to the experimental results by Tzukatani *et al.* (1997), the lines to the simulation obtained using the model (14) solved by (15) for the duplication time given there (22.3 days) and for different values of σ . From above to below $\sigma = 0.01, 0.1, 0.2$.

in cell–cell adhesion and in mechano-transduction. It seems that the link between over-expression of cadherins and growth inhibition is a characteristic shared by all types of cadherins. For instance, there are results regarding N-cadherins in the Chinese Hamster Ovary (Levenberg *et al.*, 1999), in vascular smooth muscle cells (Ugnow *et al.*, 2003) and in epithelial cells of the inner ear (Warchol, 2002), E-cadherins in human breast epithelial cells (St Croix *et al.*, 1998; Tzukatani *et al.*, 1997) and in epithelial and fibroblastoid cells (Stockinger *et al.*, 2001), VE-cadherins in vascular endothelial cells (Caveda *et al.*, 1996; Nelson & Chen, 2003) and T-cadherins in neuronal cells (Takeuchi *et al.*, 2000).

The direct involvement of cadherins was checked in several ways. For instance, Warchol (2002) found that the interaction of cells with synthetic beads coated with N-cadherin ligands spread over the substratum on which the cells live leads to growth arrest at the G1 phase of the cell cycle. Similarly, Caveda *et al.* (1996) found that coating the underlying substratum with the extracellular domain of recombinant VE-cadherin suppresses cell proliferation.

In addition to confirming that the expression of exogenous E-cadherin arrests cell growth, Stockinger *et al.* (2001) also found that a prolonged expression of E-cadherin may cause apoptosis.

Conversely, Castilla *et al.* (1999) found that the disruption of the intercellular cadherin junctions triggers the production of growth factors which contributes to the induction of proliferation.

By this mechanism, the loss of contact responsiveness can lead to deregulated growth, a phenomenon which is commonly associated with the formation of hyperplasia and malignant transformation such as gastric carcinoma (Becker *et al.*, 1994; Oda *et al.*, 1994), adenocarcinoma (Tzukita *et al.*, 1993), epithelial tumours (Cavallaro *et al.*, 2002; Christofori & Semb, 1999), colon polyps and carcinoma (Gottardi *et al.*, 2001), gynaecological cancers (Risinger *et al.*, 1994) and intimal thickening (Ugnow *et al.*, 2000) (see also the review by Hajra & Fearon, 2002). Actually, Cavallaro *et al.* (2002) considered this misperception of the presence of neighbouring cells as a fundamental milestone in the development of tumours naming it the ‘cadherin switch’ in analogy with the ‘angiogenic switch’ leading to the vascularization of tumours.

However, cadherins only represent the visible side of the process because of their transmembrane location. They are, so to speak, centre stage in the drama. In fact, there are many other members of the cast (*dramatis personae*) ‘backstage’ that can be and are involved in incorrect mechano-transduction.

First of all there are the catenins, the proteins that cadherin links to for a functional cell adhesion. The cadherin–catenin relation is well described by George & Dwivedi (2004). For instance, Stockinger *et al.* (2001) found that when epithelial cells are grown in vitro at low densities ($\leq 40\%$ confluency), they exhibited an increased β -catenin activity, which was five- to seven-fold reduced when cells reached a confluency $> 80\%$. Then in physiological conditions upon reaching confluency, the expressed cadherins sequester other catenins downregulating their activity, so that as a final result cell adhesion negatively affects cell proliferation.

Conversely, upregulation of catenins is known to induce cell proliferation by a cascade of events which will be explained in the following and is sketched in Fig. 2. In order to test the association with catenins, Chiquet *et al.* (1996) transfected Chinese Hamster Ovary cells with a cytoplasmic truncated mutant of VE-cadherin. They found that the deletion of the cytoplasmic tail does not affect the adhesive properties of VE-cadherin. However, it abolishes its growth inhibitory activity.

In more detail, as explained in Dietrich *et al.* (1997), the mechanism of contact inhibition of growth seems to be the following: the under-expression of catenins, which can be due to cell contact and over-expression of cadherins, determines the accumulation of the cyclin-dependent kinase (cdk) inhibitors p16, p21 and p27. Their over-expression inhibits entry into the S-phase causing cell cycle arrest in the G1 phase (Coats *et al.*, 1996; Kato *et al.*, 1997; Polyak *et al.*, 1994a). The increase of cellular levels of the cdk inhibitor p27 was, for instance, observed in epithelial cells growing to confluence by Polyak *et al.* (1994a) and in fibroblasts by Hengst *et al.* (1994) and Winston *et al.* (1996). The accumulation is partly due to a translational upregulation, partly due to increased half-life of p27 in response to decreased degradation by the ubiquitin-proteasome pathway (Hengst & Reed, 1996). In more detail, referring to Fig. 2, p16 blocks the activity of cdk4 by dissociating cyclin D from cdk4 and binding to cdk4. On the other hand, p27 inhibits cdk2-cyclin E directly by binding to the complex (see also Polyak *et al.*, 1994b, who also study the influence of transforming growth factor- β in the process).

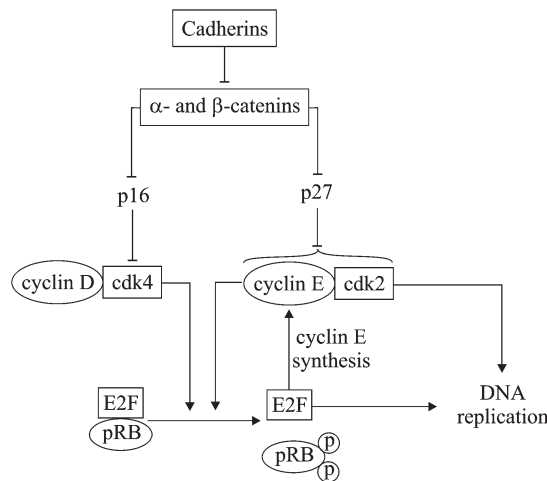


FIG. 2. Sketch of the protein cascade involved in contact inhibition of growth. cdk stands for cyclin-dependent kinase, pRB for hypophosphorylated retinoblastoma and the added p's indicate its phosphorylation. Partly modified from Dietrich *et al.* (1997). The usual arrow indicates a stimulatory action, while the other arrow indicates an inhibitory action. The curly bracket indicates that p27 binds to the cdk2/cyclin complex.

Conversely, still referring to Fig. 2, upregulation of catenins leads to the expression of cdk and then to DNA replication and mitosis.

As a further confirmation of the above cascade of events, Dietrich *et al.* (1997) found that in contact-inhibited cells the relative levels of hypophosphorylated retinoblastoma (pRB) increase.

It seems that catenin is an important crossing of several cascades leading to contact inhibition of growth. For instance, it has been recently found that α - and β -catenins bind to the tumour suppressor gene product adenomatous polyposis coli (APC) and APC mutations in the region responsible for binding to catenins are associated with the development of hyperplasia, which, as already mentioned, is an early event in tumorigenesis (Rubinfeld *et al.*, 1993; Su *et al.*, 1993).

Barker & Clevers (2000) describe how the growth suppressive-effect of E-cadherin requires the presence of its cytoplasmic β -catenin interaction domain, how catenin over-expression is related to Wnt signalling, and how catenin transcriptional activity is implicated in inducing hyper-proliferation in various tumours. In fact, it was found that increased β -catenin transcriptional activity upon over-expression of stable activated β -catenin in the intestine (Harada *et al.*, 1999) or of truncated β -catenin in the skin (Gat *et al.*, 1998) of transgenic mice leads to the development of intestinal polyp and hair tumours, respectively. β -catenin mutations have been detected in many other tumours, suggesting that a decreased cadherin/catenin association contributes to neoplastic formation (Harja & Fearon, 2002). In this paper, without being specific, we work under the assumption that something goes wrong in the cascade of events mentioned above, generating an imperfect mechano-transduction relative to the compressed state of a cell and therefore an abnormal duplication of cells.

Another issue to keep in mind in developing a macroscopic model of growing tissues is that they are not in a stress-free configuration (see Chuong & Fung, 1986; Fung, 1993; Skalak *et al.*, 1996; Taber & Humphrey, 2001). Actually, it seems that cells prefer to feel a moderate amount of stress (see Rouslahti, 1997, and references therein). A residual stress is in fact present in many cases and may be related to differential growth (Ambrosi & Guana, 2005; Ambrosi & Mollica, 2004; Humphrey & Rajagopal, 2002). For this reason, in the model we develop in this paper, we will include the possibility of duplication at moderate levels of compression.

The other main component considered in this paper is the ECM which is composed of many constituents produced by a variety of stromal cells, mainly fibroblasts. In this paper, we do not distinguish between the different types of cells in the tissue, but only distinguish between normal and abnormal cell populations and stroma. In this sense, we will generally say that the cell population produces ECM.

The percentage of cells and ECM changes considerably from tissue to tissue. For instance, the skeletal muscle collagen volume fraction is 14%, that for corneal stroma is 29% and that for subcutaneous tissue is 21% (Levick, 1987). Tendons consist of 55–70% of water and a substantial part of this is associated with proteoglycans in the ECM. Of the tendon dry weight, 66–85% is collagen (Kjaer, 2004).

Regarding the rate of production, in vitro endogenous proteins such as fibronectin and von Willebrand factor are released by endothelial cells so that a complex matrix is organized within a few hours after seeding (Chiquet *et al.*, 1996; Dejana *et al.*, 1990).

The ECM is then continuously renewed and remodelled through both the production of matrix metalloproteinases (MMPs) and the synthesis of new ECM components. In stationary physiological conditions, the remodelling of ECM is much slower than when a new tissue has to be produced. For instance, in the human lung the physiological turnover of total ECM is 10–15% per day (Johnson, 2001), which means that a complete turnover occurs in nearly a week. Our understanding of the aforementioned experimental evidence is that upon reaching confluency, the cells slow down the production of ECM components.

It is however well-known that the remodelling process is strongly affected by the stresses and strains the tissue is subject to. This is classically known for bones and teeth (Mao & Nah, 2004) and muscles (Kim *et al.*, 2002; Kjaer, 2004) and is physiologically functional because it allows for the stroma to be kept young and reactive. In fact, prolonged rest or space flight are detrimental for bones and muscles, while exercise and physical training have an opposite effect. It is common knowledge that traction is applied to heal fractures or to govern the growth of mandibular bones in children. More specifically, in bones, mechanical strain induces MMP-13 expression through the MEK-ERK signalling pathway (Yang *et al.*, 2004).

Experiments to quantify the production of ECM mainly under cyclic stretch conditions have been performed by several authors (see, for instance, Breen, 2000; Chiquet *et al.*, 2003; Coplen *et al.*, 2003; Kim *et al.*, 2002; Lee *et al.*, 2005; Yang *et al.*, 2004). Increased production of ECM, growth factor production and collagenase activity were also generated by mechanically stimulating cardiac fibroblasts (MacKenna *et al.*, 2000).

Several authors have studied the differences between the chemical and morphological characteristics of the ECM of normal tissues and that of tumours. For instance, Takeuchi *et al.* (1976) found that breast tumours presented a denser and more fibrous stroma, with variation in the content of hyaluronic acid, glycosaminoglycans, chondroitin sulfate and dermatan sulfate among other components (see also Alini & Losa, 1991; Marotta *et al.*, 1985; Paszek *et al.*, 2005; Vijayagopal *et al.*, 1998). Variations in the contents of type I collagen (Lagace *et al.*, 1985), fibronectin and laminin (Brown *et al.*, 1999; Christensen, 1992) were also observed. Christensen (1992) also argued that the above variations can be due to increased production by the tumour cells.

Increased expression of fibronectin not balanced by increased collagenase activity has also been observed elsewhere such as cardiac hypertrophy, cardiac fibrosis, liver fibrosis, pulmonary fibrosis, asthma, glomerulonephritis and colon cancer (Brewster *et al.*, 1990; Johnson, 2001; MacKenna *et al.*, 2000; Pujuguet *et al.*, 1996). In intima hyperplasia, both the amount of cells and of ECM increase (Newby & Zaltsman, 2000). Johnson (2001) argued that the alteration in the ECM components can be due to several, probably concurring, reasons: increased de novo synthesis of ECM proteins, decreased activity of its degrading enzyme (MMP), and upregulation of the tissue-specific inhibitors of metalloproteinases (TIMP).

On the other hand, excessive degradation of ECM due to excessive production of MMP-13 characterizes chronic inflammatory diseases such as osteoarthritic cartilage, rheumatoid synovium, chronic ulcer, intestinal ulcerations, periodontitis and many malignant tumours (Yang *et al.*, 2004).

The composition of ECM also changes with tumour progression and is so important in determining cell growth, differentiation and movement that Liotta & Kohn (2001) mention the possibility of finding a 'stromal therapy as a new strategy' against cancer. Referring to the changes of ECM content with tumour progression, for instance, Zhang *et al.* (2003) showed that the content of collagen fibres increased with prostate cancer grade ranging from nearly 7 to 26% of the area analysed.

3. The mathematical model

In constructing the model, we refer to the volume ratio related to both normal and abnormal cells and ECM, i.e. the fraction of volume occupied by the single constituents. We will use the term 'abnormal cells' to denote in general those cells which respond in a physiologically different way to compression and might give rise to a tumour at some later stage in their evolution. Abnormal ECM is that produced by such abnormal cells.

To be more precise

- n is the volume ratio occupied by normal cells;
- a is the volume ratio occupied by tumour cells;
- m_n is the volume ratio occupied by host ECM;
- m_a is the volume ratio occupied by ECM produced by tumour cells, which is known to be structurally and chemically different from that produced by normal cells (Alini & Losa, 1991; Lagace *et al.*, 1985; Marotta *et al.*, 1985; Takeuchi *et al.*, 1976; Vijayagopal *et al.*, 1998).

These state variables are then already dimensionless and are in the range $[0, 1]$. Considering that, as explained in Section 3, their values strongly depend on the tissue considered, throughout the paper we will work with a virtual tissue made of 50% of cells, 20% of ECM and 30% of extracellular liquid, which means that $n \approx 0.5$ and $m_n \approx 0.2$.

In addition to the volume ratios above, another fundamental state variable is

- c , the concentration of MDEs, e.g. the plasminogen activators or MMPs.

Unlike the cells and ECM, MDEs are considered as macromolecules which diffuse in the extracellular liquid without occupying space.

In the following, an important role will be played by the overall volume ratio occupied by cells and ECM

$$\psi = n + a + m_n + m_a \quad (1)$$

and by its relation to the stress. In using (1), it is assumed for simplicity that all cells and matrix constituents of any origin contribute equally to the sensing of stress by the cells. A discussion of this hypothesis is given in De Angelis & Preziosi (2000) where alternatives are also proposed.

Several forms of stress–volume ratio relations can be taken into account (Ambrosi & Preziosi, 2002; Byrne & Preziosi, 2004). The simplest features characterizing stress are that it vanishes below a value ψ_0 , which we denote as the stress-free volume ratio and corresponds to the confluent volume ratio, is increasing for $\psi > \psi_0$ and tends to infinity at $\psi = 1$, e.g.

$$\Sigma(\psi) = E(1 - \psi_0) \left(\frac{\psi - \psi_0}{1 - \psi} \right)_+, \quad (2)$$

where $(f)_+$ denotes the positive part of f and E is the value of the derivative in $\psi = \psi_0$, a sort of Young modulus for moderate compressions. However, since usually $\psi \approx \psi_0$, then one can also approximate (2) with $\Sigma(\psi) = E(\psi - \psi_0)_+$. We explicitly mention that the constitutive relation (2) is not able to describe the response to tensile stresses and adhesion-like behaviours.

The mass balance equations for the two types of cells in our model can then be written as

$$\begin{aligned} \frac{\partial n}{\partial t} + \nabla \cdot (n\mathbf{v}_n) &= \Gamma_n n - \delta_n n, \\ \frac{\partial a}{\partial t} + \nabla \cdot (a\mathbf{v}_a) &= \Gamma_a a - \delta_a a, \end{aligned} \quad (3)$$

where we will consider the death coefficient to be constant, though apoptosis may be influenced by the stress level (Helmlinger *et al.*, 1997).

Referring to the phenomenological discussion presented in Section 2, we assume that cells replicate if they feel there is sufficient space to do this, i.e. if they can sense a sustainable level of compression

($\psi \leq \psi_n$), or equivalently stress ($\Sigma \leq \Sigma_n$ with $\Sigma_n = \Sigma(\psi_n)$). Our assumption is that the threshold value ψ_n is nearly equal to the stress-free value ψ_0 for normal cells and slightly larger for abnormal ones ($\psi_a > \psi_n$). Actually, growth might even be independent of ψ (or equivalently $\psi_a \geq 1$), meaning that the cells are completely insensitive to mechanical cues and continue replicating independently of the compression level.

To describe mathematically what we have stated above, we choose the simplest possible switch mechanism depending on compression for the growth coefficient, i.e.

$$\begin{aligned} \Gamma_n(\psi) &= \gamma_n H_\sigma(\psi - \psi_n), \\ \Gamma_a(\psi) &= \gamma_a H_\sigma(\psi - \psi_a), \end{aligned} \quad (4)$$

where H_σ is a monotonic mollifier of the step function with the properties that it is a continuous function with $H_\sigma(\phi) = 1$ if $\phi \leq 0$ and $H_\sigma(\phi) = 0$ if $\phi \geq \sigma$. In particular, we will choose

$$H_\sigma(\phi) = \begin{cases} 1, & \text{if } \phi \leq 0, \\ 0, & \text{if } \phi > \sigma, \\ 1 - \frac{\phi}{\sigma}, & \text{otherwise.} \end{cases} \quad (5)$$

From the biological point of view, σ measures the mechano-transduction sensitivity. We explicitly note that for $\psi_n > \psi_0$, H_σ can be equivalently defined as a switch mechanism based on the stress level. In fact, in the following we will often make use of the correspondence between compression and stress.

Of course, the growth term depends on other quantities, e.g. the amount of nutrient and growth factors. In this paper, however, we want to focus on the possible role of compression and stress on tumour invasion and therefore assume that all nutrients are abundantly supplied and neglect the effect of growth factors. In doing this, we tacitly assume that all the constituents necessary for the cell to grow and undergo mitosis can be found in the extracellular liquid that is a passive constituent in the global mass balance equation. Actually, in the following examples we will consider the case in which $\gamma_n = \gamma_a$ meaning that the reproduction rate is the same and only the stress perception is different $\psi_n < \psi_a$. For the same reason, we will not consider other mechanisms such as the reduced viability of normal tissue in a low-pH environment (Gatenby & Gawlinski, 1996).

Regarding the motion of normal cells, observing that they move in the intricate network formed by the ECM (see the supplementary material for Wolf *et al.*, 2003), we model the ensemble of cells as a granular material in a porous medium. Therefore, following Ambrosi & Preziosi (2002) and Preziosi & Graziano (2003), we can write the momentum equation for the cells

$$\rho n \left(\frac{\partial \mathbf{v}_n}{\partial t} + \mathbf{v}_n \cdot \nabla \mathbf{v}_n \right) = \nabla \cdot \mathbf{T}_n + \mathbf{m}_n, \quad (6)$$

which, neglecting inertia and introducing the usual porous media assumptions on the stress tensor \mathbf{T}_n and on the interaction force \mathbf{m}_n , leads to

$$\mathbf{v}_n = -K \nabla \Sigma(\psi), \quad (7)$$

where K is related to the permeability of the medium.

In addition, we explicitly mention that in many of the cases that we will deal with in the following sections, the exact form of the stress-volume ratio relation is of minor importance.

The coefficient K is not only related to the pure permeability of the porous medium but also related to the ability of cells to move through the ECM network. It will in fact depend on the amount of ECM

for two main reasons: space occupation and the presence of adhesion sites. Actually, increased motility is observed for regions with higher concentration of ECM. However, it decreases again for large values of ECM concentration because of the increase of adhesive sites (see Hillen, 2002; Painter & Hillen, 2002; Perumpanani & Byrne, 1996).

The closure (7), phenomenologically introduced in Bertsch *et al.* (1985) and Gurtin & McCamy (1977), has the effect to push cells away from the overcrowded region where they feel pressed, until they reach, if possible, a stress-free configuration. In mathematical terms, this means that the motion is down the stress gradient, or the overall density gradient.

In the following, we will assume that normal and abnormal cells behave in the same way under compression, i.e.

$$\mathbf{v} \equiv \mathbf{v}_a = \mathbf{v}_n = -K \Sigma'(\psi) \nabla \psi, \quad (8)$$

so that (3) can be re-written as

$$\begin{aligned} \frac{\partial n}{\partial t} &= \nabla \cdot [nK \Sigma'(\psi) \nabla \psi] + [\gamma_n H_\sigma(\psi - \psi_n) - \delta_n]n, \\ \frac{\partial a}{\partial t} &= \nabla \cdot [aK \Sigma'(\psi) \nabla \psi] + [\gamma_a H_\sigma(\psi - \psi_a) - \delta_a]a. \end{aligned} \quad (9)$$

The ECM, which as has already been mentioned, is an intricate network of fibrous material made of many macromolecules, including fibronectin, laminin and collagen, is assumed here to be rigid. However, we take into account that it can be degraded by MDEs (Chambers & Matrisian, 1997; Matrisian, 1992; Parson *et al.*, 1997; Stetler-Stevenson *et al.*, 1996). Active MDEs are produced (or activated) by the cells, diffuse throughout the tissue and undergo some form of decay, either passive or active. The equation governing the evolution of MDE concentration is therefore given by

$$\frac{\partial c}{\partial t} = \kappa \nabla^2 c + P(n, a, c, \Sigma) - D(c, m)c. \quad (10)$$

The functions P and D model the production of active MDEs by normal and abnormal cells and their decay, respectively. We will take $D = 1/\tau$ to be a constant (other functional forms for D have also been investigated, see Anderson *et al.*, 2000, for details) and P to be of the following form:

$$P = \pi_n(\Sigma)n + \pi_a(\Sigma)a. \quad (11)$$

In the following, we will consider in more detail the case in which π_n and π_a are constant.

Upon contact, MDEs degrade the ECM produced by the cells in a stress-dependent way (see, for instance, MacKenna *et al.*, 2000; Yang *et al.*, 2004). Hence, the degradation process can be modelled by the following simple equations:

$$\begin{aligned} \frac{\partial m_n}{\partial t} &= \mu_n(\Sigma)n - \nu c m_n, \\ \frac{\partial m_a}{\partial t} &= \mu_a(\Sigma)a - \nu c m_a, \end{aligned} \quad (12)$$

where ν is a positive constant and μ_n and μ_a depend on whether the cells are in a confluent situation or not, or on the fact that the tissue is subject to strain. In most of the simulations to follow we will take

them as constant. The study of the case in which π_i and μ_i ($i = n, a$) depend on the level of stress, which as already mentioned is of great interest in tissue remodelling, is left as a future development when more quantitative experimental data on these dependencies are available.

In (12), we distinguish the production of ECM by normal and abnormal cells, because as explained in Section 2, they are morphologically and chemically different. In addition, the production rates can be different, while it is known that the degradation by MDEs does not distinguish between the two types of ECM.

Summarizing, the complete system of equations is

$$\left\{ \begin{array}{l} \frac{\partial n}{\partial t} = \overbrace{\nabla \cdot [nK \Sigma'(\psi) \nabla \psi]}^{\text{response to compression}} + \overbrace{\gamma_n H_\sigma(\psi - \psi_n)n}^{\text{growth}} - \overbrace{\delta_n n}^{\text{apoptosis}} , \\ \frac{\partial a}{\partial t} = \overbrace{\nabla \cdot [aK \Sigma'(\psi) \nabla \psi]}^{\text{response to compression}} + \overbrace{\gamma_a H_\sigma(\psi - \psi_a)a}^{\text{growth}} - \overbrace{\delta_a a}^{\text{apoptosis}} , \\ \frac{\partial m_n}{\partial t} = \overbrace{\mu_n(\Sigma(\psi))n}^{\text{host production}} - \overbrace{vcm_n}^{\text{degradation}} , \\ \frac{\partial m_a}{\partial t} = \overbrace{\mu_a(\Sigma(\psi))a}^{\text{tumour production}} - \overbrace{vcm_a}^{\text{degradation}} , \\ \frac{\partial c}{\partial t} = \overbrace{\kappa \nabla^2 c}^{\text{diffusion}} + \overbrace{\pi_n(\Sigma(\psi))n}^{\text{host production}} + \overbrace{\pi_a(\Sigma(\psi))a}^{\text{tumour production}} - \overbrace{\frac{c}{\tau}}^{\text{decay}} , \end{array} \right. \quad (13)$$

where H_σ and Σ are defined by (2) and (5), respectively, and we recall that $\psi = n + a + m_n + m_a$ is the percentage of volume not occupied by the extracellular liquid.

Of course System (13) requires to be closed with appropriate initial and boundary conditions, which depend on the specific problem of interest. However, in general zero-flux boundary conditions will be imposed, while, as far as the initial conditions are concerned, a given distribution of abnormal cells $a(t = 0, x) = a_0(x)$ will be assumed to replace the normal cells starting with $m_n(t = 0, x) = 0$. The other variables will have initial conditions as the stationary values described in Section 4.

4. Qualitative behaviour of spatially independent solutions

Before proceeding with the numerical simulation of the full model, we consider some simple cases in order to understand and describe the main features of the model presented above in the spatially homogeneous case. This allows us to give the analytical form of some significant spatially independent solutions.

4.1 Generation of confluent distribution in vitro

In order to compare the results of the model with some in vitro experiments such as the one shown in Fig. 1 and to identify some of the parameters of the model, we start with the easiest situation in which a few cells are uniformly seeded on a suitable substratum in a Petri dish and, duplicating on the surface of the substratum, grow to confluence with no production of ECM. This situation is modelled by assuming spatial homogeneity and $m_n(t = 0, x) = m_a(t = 0, x) = a(t = 0, x) = 0$, so that one can restrict to

consider only

$$\frac{\partial n}{\partial t} = [\gamma_n H_\sigma (n - \psi_n) - \delta_n]n, \quad (14)$$

which for H_σ given by (5) can be easily solved by

$$n = \begin{cases} n_0 \exp[(\gamma_n - \delta_n)t], & \text{for } t \leq t_0, \\ \frac{\psi_n + (1 - \frac{\delta_n}{\gamma_n})\sigma}{1 + (1 - \frac{\delta_n}{\gamma_n})\frac{\sigma}{\psi_n} \exp[-(\gamma_n - \delta_n + \gamma_n \frac{\psi_n}{\sigma})(t - t_0)]}, & \text{for } t > t_0, \end{cases} \quad (15)$$

where $n_0 < \psi_n$ is the initial volume ratio of cells and

$$t_0 = \frac{1}{\gamma_n - \delta_n} \ln \frac{\psi_n}{\sigma} \quad (16)$$

is when confluency starts. The solution will tend towards the volume ratio

$$n_\infty = \hat{\psi}_n := \psi_n + \left(1 - \frac{\delta_n}{\gamma_n}\right)\sigma, \quad (17)$$

which, even if $\psi_n = \psi_0$, implies a small compression of the tissue $\Sigma \approx E(1 - \frac{\delta_n}{\gamma_n})\sigma$. In Fig. 1, we compare the results of the experiments by Tzukatani *et al.* (1997) and (15) for the doubling time $T = 22.3$ h ($\gamma_n = 0.746$ per days) measured therein and for different values of σ assuming that the stationary cell density measured there ($N = 69000$ cells over a 35-mm dish) correspond to $\psi_n + \sigma = 0.7$. In the experiments by Tzukatani *et al.* (1997), an unexplained decrease of the cell population in the first day is measured. For this reason, the cell density measured 48 h after seeding was taken as initial condition for (14) ($N_0 = 45000$ cells over a 35-mm dish, corresponding to $n_0 = 0.495$). From the simulation it appears that $\sigma \approx 0.2$. For the sake of completeness in Table 1 we report the growth rates measured for different cell types.

4.2 Generation of normal tissue

We consider now how the model describes the re-construction of a normal tissue starting from a few cells $n_0 < \psi_n$ constantly distributed in space. Possibly, in the same region there might be a small constant

TABLE 1 *Growth rates for different cell types. CHO stands for Chinese Hamster Ovary cells, HBE for Human Breast Epithelial cells BPAEC for Bovin Pulmonary Artery Endothelial Cells*

Cell type	Duplication time (h)	Growth rate (per days)	Source
CHO	18–19	0.86–0.91	Levenberg <i>et al.</i> (1999)
CHO	14–21	0.8–1.2	Caveda <i>et al.</i> (1996)
L929	10–12	1.4–1.6	Caveda <i>et al.</i> (1996)
Fibroblast	16.2	1.03	Stockinger <i>et al.</i> (2001)
HBE	22.3	0.746	Tzukatani <i>et al.</i> (1997)
BPAEC	38–44	0.38–0.44	Nelson & Chen (2003)

amount of ECM, so that at least initially $n + m_n < \psi_n$. In this case, the system initially reduces to

$$\begin{cases} \frac{\partial n}{\partial t} = [\gamma_n H_\sigma (n + m_n - \psi_n) - \delta_n]n, \\ \frac{\partial m_n}{\partial t} = \mu_n n - \nu c m_n, \\ \frac{\partial c}{\partial t} = \pi_n n - \frac{c}{\tau}. \end{cases} \quad (18)$$

Denoting by $\hat{\psi}_n \in [\psi_n, \psi_n + \sigma]$ the value such that $H_\sigma(\hat{\psi}_n - \psi_n) = \delta_n/\gamma_n$ (given by (17) if H_σ is given by (5)), the steady state can easily be written down

$$n = \hat{\psi}_n - M_n, \quad m_n = M_n, \quad c = \pi_n \tau (\hat{\psi}_n - M_n), \quad (19)$$

where

$$M_n = \frac{\mu_n}{\nu \pi_n \tau}. \quad (20)$$

In particular, the ratio of cells to ECM is

$$\eta := \frac{n}{m_n} = \frac{\nu \pi_n \tau}{\mu_n} \hat{\psi}_n - 1, \quad (21)$$

which is an easily measurable value. We observe that M_n has to be smaller than 1, otherwise the production of ECM would never be balanced by degradation leading to the formation of a tissue only made of ECM with no cells.

In order to make explicit the analytical form of the solution, the discussion proposed in this and in the following subsections holds in the limit $\sigma \rightarrow 0$, from the biological point of view, corresponding to a sharp transition between the quiescent and the proliferative states. In fact, for instance, in this case (18) can be solved sequentially (i.e. having integrated the equation for n , one can integrate the equation for c and finally the equation for m_n). However, for the sake of simplicity, considering that, as will be explained in the discussion of the parameter estimates below (see also Hutton *et al.*, 1998; Ottl *et al.*, 2000; Tzukatani *et al.*, 1997), MDE production and degradation is much faster than cell duplication, and we only write here the solution obtained under the assumption that $c(t) \approx \pi_n \tau n(t)$. In this case, the solution can be written as

$$n = n_0 \exp[(\gamma_n - \delta_n)t], \quad (22)$$

$$m_n = M_n + (m_0 - M_n) \exp\left[\frac{\nu \pi_n \tau n_0}{\gamma_n - \delta_n} (1 - e^{(\gamma_n - \delta_n)t})\right], \quad (23)$$

until $m_n(t) + n(t)$ reaches ψ_n . From this instant, which will be denoted by \hat{t} and can be determined through an implicit relation, $n(t) = \psi_n - m_n(t)$ where $m_n(t)$ satisfies

$$\frac{\partial m_n}{\partial t} = \mu_n \left(1 - \frac{m_n}{M_n}\right) (\psi_n - m_n), \quad (24)$$

starting with the ECM volume ratio $\hat{m}_n = \psi_n - n_0 \exp[(\gamma_n - \delta_n)\hat{t}]$ reached at $t = \hat{t}$, which is solved for $t > \hat{t}$ by

$$m_n = \frac{M_n + \psi_n \frac{\hat{m}_n - M_n}{\psi_n - \hat{m}_n} \exp\left[-\mu_n \left(\frac{\psi_n}{M_n} - 1\right)(t - \hat{t})\right]}{1 + \frac{\hat{m}_n - M_n}{\psi_n - \hat{m}_n} \exp\left[-\mu_n \left(\frac{\psi_n}{M_n} - 1\right)(t - \hat{t})\right]}. \quad (25)$$

4.2.1 Parameter estimation. In addition to the growth rates already summarized in Table 1, we will quantify here, as far as possible, the parameters involved before proceeding with the simulation. From the observation that *in vitro* cells organize an ECM within a few hours after seeding (Chiquet *et al.*, 1996; Dejana *et al.*, 1990) we infer that $\mu_n \approx 10$ per days. On the other hand, from the fact that 10–15% of the lung ECM is renewed every day (Johnson, 2001) we can evaluate $\mu_n \approx 0.1$ per days. This suggests that the ECM production decreases while cells achieve a confluent configuration. In most of the following simulations we will use a value near the lower bound because we usually deal with ECM remodelling.

The other parameters directly measured are the diffusion coefficient of MDEs, specifically MMP2, $\kappa = 0.85 \times 10^{-6} \text{ cm}^2/\text{s}$ (Berk *et al.*, 1993) and the MDE half-life, which of course depends on the specific enzyme and may vary anywhere in the range 10^{-4} –0.13 h (Hutton *et al.*, 1998; Ottl *et al.*, 2000). This means that the range of action $\zeta = \sqrt{\kappa\tau}$ of the various MDEs varies from a few microns, i.e. a very local action, to a few hundreds of microns, which is one order of magnitude less than the typical size of an avascular tumour at its diffusion-limited steady state. However, to consider the entire chain of reactions we will take a value near the upper limit for τ .

Karagiannis & Popel (2004) estimated MMP production rate to be approximately 1800000 molecules/cell h, which leads to a production rate π_n of the order of 10^6 – 10^7 molecules/ μm^3 days. Since the steady value foreseen by the model is $c_\infty = \pi_n\tau n_\infty$, then we assume that the concentration of MDEs is about 15000 molecules/ μm^3 . Finally, since we are working with an ideal tissue made of 20% ECM, then $v \approx 10^{-5} \mu\text{m}^3/\text{molecules days}$.

Finally, the apoptotic rate δ_n , which has to be less than γ_n , is assumed to be 0.1 per days.

The corresponding values for abnormal tissue are a slight modification of the values above, and these will be quantified in the discussion and the simulation to follow.

The above parameters' estimates have been used in the simulation given in Fig. 3. The solution presents an overshooting because of the fact that the volume ratios of ECM and cells do not reach the

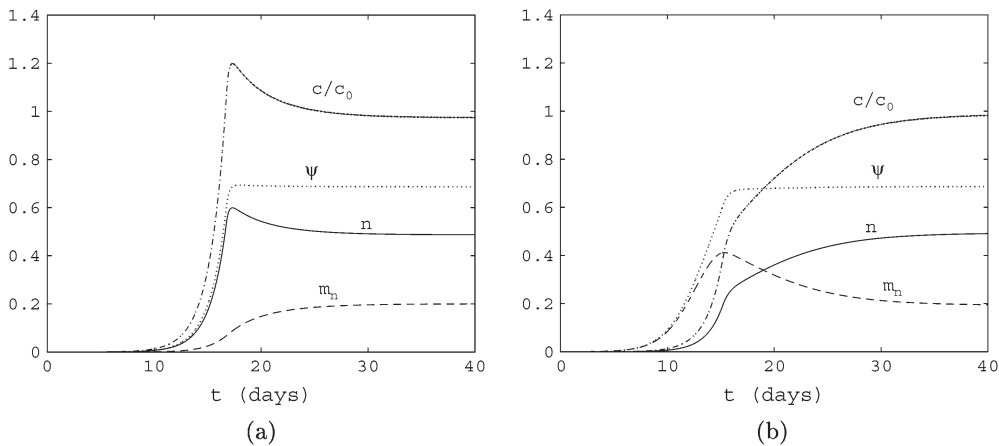


FIG. 3. Generation of a normal tissue in the case of a constant slow production of ECM (a) and of a density-dependent production approximating a fast production at early times and a slow stationary production at confluency (b). The values of the parameters are $\gamma_n = 0.746$ per days, $\delta_n = 0.1$ per days, $\pi_n/c_0 = 400$ per days, $\nu c_0 = 0.25$ per days, $\tau = 0.005$ days, $\psi_n = 0.6$, $\sigma = 0.1$ where c_0 is the reference concentration value 15000 molecules/ μm^2 . In (a) $\mu_n = 0.1$ per days. From bottom to top, volume ratio of ECM, cells and overall volume ratio and normalized MDE concentration. In (b) $\mu_n = 10(1 - \psi)^4$ per days leading to a faster initial production of ECM and a temporary hyper-production of ECM.

stationary value at the same time. In fact, the production rates are different. In particular, in Fig. 3(a) the production of ECM is quite slow because right from the beginning, the ECM production rate is of the order of the value estimated at the steady state. Instead, in Fig. 3(b) $\mu_n = 10(1 - \psi)^4$ per days is chosen as an example of a decreasing production of ECM. The difference with Fig. 3(a) consists in a larger production of ECM at earlier times which determines an initial overproduction of ECM occupying up to twice the volume in the final configuration. The extra ECM is then progressively degraded by the MDEs to reach the stationary value $\mu(\hat{\psi}_n)/\nu\pi_n\tau$.

4.3 Development of abnormal tissue

Assume now that by a clonal differentiation, some abnormal cells appear in the fully developed tissue. We first assume that they appear constantly throughout the domain, so that the system now becomes

$$\begin{cases} \frac{\partial n}{\partial t} = [\gamma_n H_\sigma(\psi - \psi_n) - \delta_n]n, \\ \frac{\partial a}{\partial t} = [\gamma_a H_\sigma(\psi - \psi_a) - \delta_a]a, \\ \frac{\partial m_n}{\partial t} = \mu_n n - \nu c m_n, \\ \frac{\partial m_a}{\partial t} = \mu_a a - \nu c m_a, \\ \frac{\partial c}{\partial t} = \pi_n n + \pi_a a - \frac{c}{\tau}, \end{cases} \quad (26)$$

a natural modification of the stationary solution (19).

We take as initial conditions

$$\begin{aligned} n(t = 0, x) &= \psi_n - a_0 - M_n, \\ a(t = 0, x) &= a_0, \\ m_n(t = 0, x) &= M_n, \\ m_a(t = 0, x) &= 0, \\ c(t = 0, x) &= \pi_n \tau (\psi_n - M_n), \end{aligned} \quad (27)$$

where M_n is defined in (20).

We start by assuming that the rates of production of ECM and MDEs are equal ($\mu_n = \mu_a$ and $\pi_n = \pi_a$). The equations for the ECM can then be added to give

$$\frac{\partial m}{\partial t} = \mu_n(n + a) - \nu c m, \quad (28)$$

where $m = m_n + m_a$.

Observing, as in Section 4.2.1, that the timescales of production and degradation of MDEs are much faster than the time of reproduction of the cells and of production of the more complex ECM (minutes with respect to hours or days, see Table 2), we can state that the concentration of MDEs quickly reaches the value

$$c = \pi_n \tau (n + a), \quad (29)$$

which can be substituted back in (28) to obtain

$$\frac{\partial m}{\partial t} = \mu_n \left(1 - \frac{\pi_n \nu \tau}{\mu_n} m \right) (n + a). \quad (30)$$

TABLE 2 *Parameter estimate*

Parameter	Estimated value	Source
δ	0.1 per days	See text
μ_n	0.1–10 per days	Chiquet <i>et al.</i> (1996), Dejana <i>et al.</i> (1990) and Johnson (2001)
ν	$10^{-5} \mu\text{m}^2/\text{molecules days}$	See text
π_n	$10^6\text{--}10^7 \text{ molecules}/\mu\text{m}^2\text{days}$	See text
τ	$\frac{1}{6} \text{ s--}0.13 \text{ h}$	Hutton <i>et al.</i> (1998) and Ottl <i>et al.</i> (2000)
κ	$0.85 \times 10^{-6} \text{ cm}^2/\text{s}$	Berk <i>et al.</i> (1993)

Therefore, from the initial conditions, one has that

$$m(t) = m_n(t) + m_a(t) = M_n, \quad (31)$$

remains constant which, in particular, implies $\psi(t) = n(t) + a(t) + M_n$ at all times.

We can then focus on the first two equations to explain first how the abnormal cells progressively replace the normal ones in the tissue.

Since $\psi_n < \psi_a$, initially the abnormal tissue continues its proliferation and

$$a(t) = a_0 e^{(\gamma_a - \delta_a)t}. \quad (32)$$

It should be noted that in this initial phase, the amount of normal cells then adjusts so that $n(t) + a(t) + M_n = \psi_n$. In fact, a solution similar to (32) would hold for the normal tissue if $\psi \leq \psi_n$, but this contradicts the fact that all populations grow starting from an initial condition with $\psi = \psi_n$. Hence, throughout the process $\psi \geq \psi_n$.

On the other hand, whenever $\psi > \psi_n$ normal cells would decay exponentially. So, n cannot overcome the value of the volume ratio corresponding to the natural configuration (actually, when using a mollifier of width σ , the volume ratio = $\psi_n + \sigma$).

Similarly to what we have just discussed, initially we cannot have $\psi > \psi_n$ positive, otherwise the total volume ratio would be

$$\psi(t) = (\psi_n - M_n - a_0)e^{-\delta_n t} + a_0 e^{(\gamma_a - \delta_a)t} + M_n, \quad (33)$$

which is smaller than ψ_n for small values of a_0 . This simply corresponds to the fact that in a small interval of time Δt , an amount $a_0(\gamma_a - \delta_a)\Delta t$ of abnormal cells is produced and a larger amount $(\psi_n - M_n - a_0)\delta_n\Delta t$ of normal cells dies (if $a_0(\psi_n - M_n) < \delta_n/(\gamma_a - \delta_a + \delta_n)$). What happens is that an amount of normal cells is still produced to adjust ψ to the maximum stress-free value. Therefore,

$$n(t) = \psi_n - M_n - a_0 e^{(\gamma_a - \delta_a)t}, \quad (34)$$

so that $n + a = \psi - M_n$ is always equal to ψ_n . This situation continues until the rate at which abnormal cells are produced equals the rate at which normal cells would normally die. Denoting by A the amount of abnormal cells at this instant and observing that the amount of normal cells is $\psi_n - M_n - A$, from the first two equations in (26), the above situation occurs when

$$A(\gamma_a - \delta_a)\Delta t = (\psi_n - M_n - A)\delta_n\Delta t. \quad (35)$$

Hence,

$$A = \frac{(\psi_n - M_n)\delta_n}{\gamma_a - \delta_a + \delta_n}, \quad (36)$$

and since $A = a(t_1) = a_0 e^{(\gamma_a - \delta_a)t_1}$, the instant is given by

$$t_1 = \frac{1}{\gamma_a - \delta_a} \log \left(\frac{\delta_n}{\gamma_a - \delta_a + \delta_n} \frac{\psi_n - M_n}{a_0} \right), \quad (37)$$

and

$$n(t_1) = \frac{\gamma_a - \delta_a}{\gamma_a - \delta_a + \delta_n} (\psi_0 - M_n). \quad (38)$$

Referring to Fig. 4(a), the period $[0, t_1]$ is characterized by a simple partial replacement of normal cells by abnormal cells, keeping the total volume ratio constantly equal to ψ_n (therefore in this phase there is no compression). For this reason we will call this phase ‘relaxed replacement’ or Phase I.

Note that if the only difference in the behaviour of normal and abnormal cells is in the sensing of stress by the cells, i.e. $\gamma_n = \gamma_a = \gamma$ and $\delta_n = \delta_a = \delta$, then

$$t_1 = \frac{1}{\gamma - \delta} \log \left(\frac{\delta}{\gamma} \frac{\psi_n - M_n}{a_0} \right), \quad (39)$$

and

$$n(t_1) = \left(1 - \frac{\delta}{\gamma}\right) (\psi_n - M_n), \quad a(t_1) = \frac{\delta}{\gamma} (\psi_n - M_n). \quad (40)$$

After the end of the relaxed replacement phase, the amount of abnormal cells produced is larger than the amount of normal cells dying, so the total volume ratio begins to exceed the stress-free value ψ_n . Abnormal cells continue to grow satisfying (32) while

$$n(t) = \frac{\gamma_a - \delta_a}{\gamma_a - \delta_a + \delta_n} \left(\frac{\delta_n}{\gamma_a - \delta_a + \delta_n} \frac{\psi_n - M_n}{a_0} \right)^{\frac{\delta_n}{\gamma_a - \delta_a}} (\psi_n - M_n) e^{-\delta_n t}. \quad (41)$$

This change can be identified in Fig. 4 by the change of convexity of the curve related to n . In this second phase,

$$\psi(t) = \frac{\gamma_a - \delta_a}{\gamma_a - \delta_a + \delta_n} \left(\frac{\delta_n}{\gamma_a - \delta_a + \delta_n} \frac{\psi_n - M_n}{a_0} \right)^{\frac{\delta_n}{\gamma_a - \delta_a}} (\psi_n - M_n) e^{-\delta_n t} + a_0 e^{(\gamma_a - \delta_a)t} + M_n > \psi_n, \quad (42)$$

giving rise to a progressive compression of the tissue as in the last cartoon in Fig. 4(a). For this reason, this phase is called the ‘compressing phase’ or Phase II. This situation continues till ψ reaches the value ψ_a .

After this a situation in some sense reciprocal to the relaxed replacement phase occurs, with a decreasing n satisfying (41) while abnormal cells replicate to replace the space left available by n , i.e.

$$a(t) = \psi_a - n(t), \quad (43)$$

so that the total volume ratio is always equal to ψ_a . This phase then will be called the ‘over-compressed replacement’ or Phase III. Eventually, abnormal cells completely replace normal cells.

Focussing now on the two types of ECMs, using (32), (34), (41) and (43) it is possible to solve the related ODEs explicitly. As expected, it is found that m_a increases and m_n correspondingly decreases to zero keeping the sum constant. The concentration of MDEs starting from a value controlling the production of normal ECM goes to a value which controls the production of tumour-produced ECM. This value is slightly higher because of the increased density of cells. So, eventually the abnormal stroma completely replaces the normal one, as shown in Figs 4 and 5.

In conclusion, a very small incorrect sensing of the state of compression by the cells will lead to a clonal advantage and a complete replacement of normal tissue by hyperplastic over-compressed tissue with $a = \psi_a - M_n$ and $m_a = M_n$.

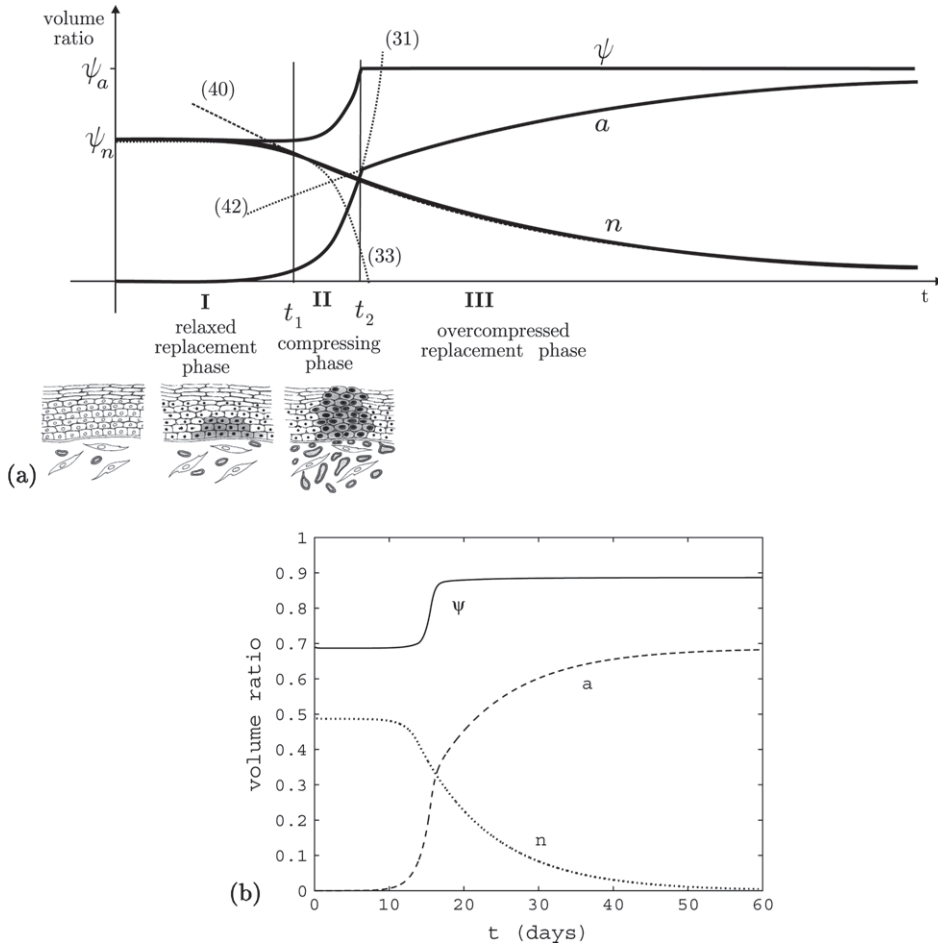


FIG. 4. Replacement of normal cells by abnormal cells. In (a) the definition of the phases of growth and of the analytical solutions (dotted lines) is given in the limit case $\sigma = 0$ (and $m = 0$). The cartoons at the bottom of the figure sketch the phenomenological evolution (modified from Hanahan & Folkman, 1996). In (b) simulation for $\gamma_n = \gamma_a = 1$ per days, $\delta_n = \delta_a = 0.1$ per days, $a_0 = 0.00001$, $\psi_n = 0.6$, $\psi_a = 0.8$ and $\sigma = 0.1$. All values of the parameters are equal for the two populations except the stress-related thresholds for stopping replication. This implies that the volume ratio of ECM is always equal to the initial value $m_0 = 0.2$.

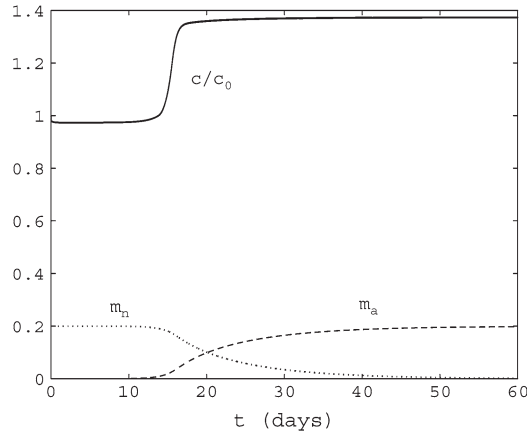


FIG. 5. Replacement of ECM in the tissue. Normal ECM (full) is replaced by the ECM produced by the tumour cells (dashed). Their sum is always equal to 0.2. The dotted line describes the evolution of the concentration of MDEs (normalized with respect to the reference value $c_0 = 15000$ molecules/ μm^2) that starting from a value controlling the production of normal ECM goes to a value which controls the production of tumour-produced ECM. This value is higher because the density of cells is higher (see Fig.4(b)). The values of the parameters are as in Fig. 4 and $\mu_n = \mu_a = 0.1$ per days, $\pi_n/c_0 = \pi_a/c_0 = 400$ per days, $\nu c_0 = 0.25$ per days, $\tau = 0.005$ days.

4.4 Hyper- and hypo-production of ECM and MDEs

If the production of MDEs and/or ECM by the two types of cells occurs at different rates, i.e. $\pi_n \neq \pi_a$ and/or $\mu_n \neq \mu_a$, the steady state derived in Section 4.3 then becomes

$$n = m_n = 0, \quad a = \psi_a - M_a, \quad m_a = M_a, \quad c = \pi_a \tau (\psi_a - M_a), \quad (44)$$

where

$$M_a = \frac{\mu_a}{\nu \pi_a \tau}, \quad (45)$$

i.e. as in the case of equal production of ECM and MDEs by the two types of cells, the abnormal ECM produced by the abnormal tissue will completely replace the ECM produced by the normal cells. However, in this case the final ECM volume ratio might be different from the initial one and depends on the dimensionless parameter M_a .

In particular, the ratio of the content of ECM produced in the normal and in the abnormal case is

$$\frac{m_n}{m_a} = \frac{\mu_n \pi_a}{\mu_a \pi_n},$$

which is in line with Johnson's observation (Johnson, 2001) that the alteration in the ECM components can be due to both increased de novo synthesis of ECM proteins, and decreased activity of its degrading enzyme (MMP), probably together with an upregulation of the TIMP.

In Fig. 6, we consider the pathological situation in which the abnormal cells are also characterized by a doubling or halving of the rate of production of ECM (Fig. 6(b) and 6(a), respectively) or of MDEs (Fig. 6(c) and 6(d), respectively).

A hyper-production of ECM (Fig. 6(b)) and a hypo-production of MDEs (Fig. 6(d)) lead to a situation where there is more fibrotic tissue with a ratio of cells versus ECM content nearly equal to 1.22, compared with 2.45 in the physiological situation. On the other hand, a hypo-production of ECM

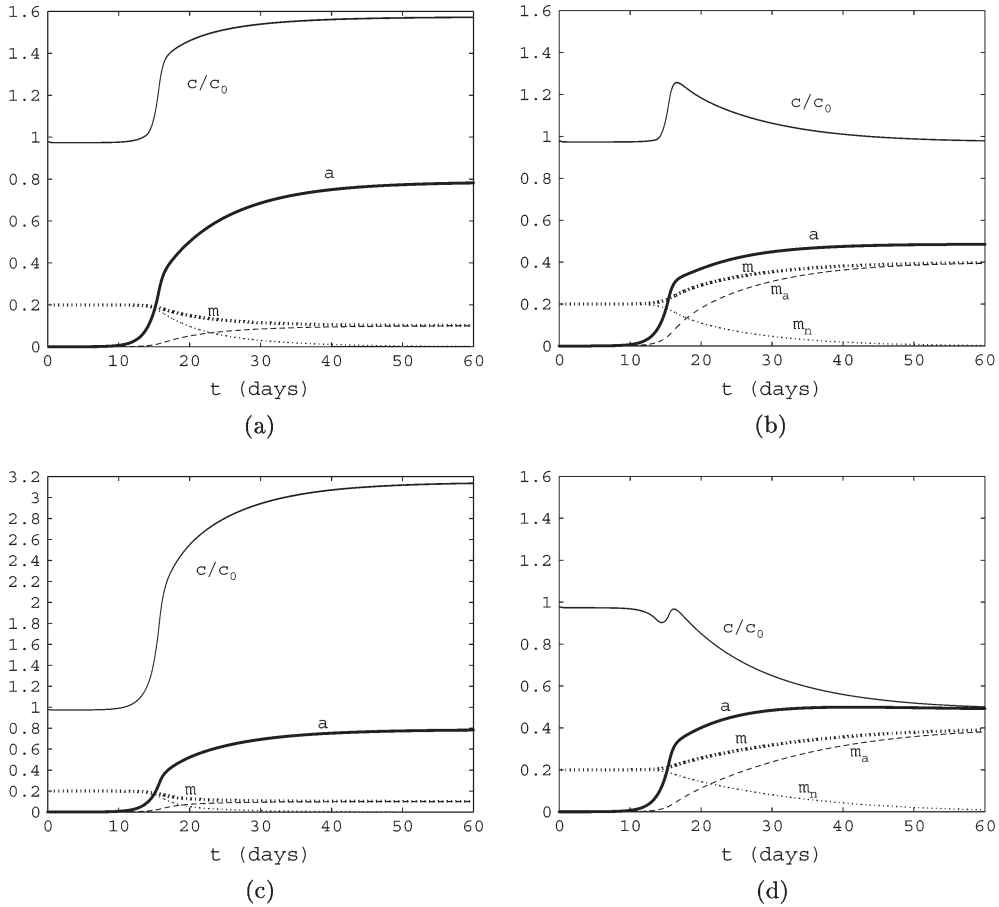


FIG. 6. Hyperplasia with complete replacement of the tissue in the case of a pathological production of ECM and/or MDEs with respect to the situation in Figs 4(b) and 5. In all figures, full lines refer to scaled concentration of MDEs, thick full lines to the volume ratio of abnormal cells, dashed lines to the volume ratio of ECM produced by abnormal cells, dotted lines to that produced by normal cells and thick dotted lines to their sum. The evolution of ψ and n is very similar to the one in Fig. 4(b) and is almost unaffected by the parameters used. In (a) $\mu_a = 0.05$ per days, in (b) $\mu_a = 0.2$ per days, in (c) $\pi_a/c_0 = 800$ per days and in (d) $\pi_a/c_0 = 200$ per days with all other parameters as in Figs 4 and 5.

(Fig. 6(a)) and a hyper-production of MDEs (Fig. 6(c)) lead to a situation where the tissue is characterized by a ratio of cells versus ECM content nearly equal to 7.86. In all cases, the evolution of ψ and n is very similar to that in Fig. 4(b).

If, for instance, both production rates increase while keeping the same ratio, then the final configuration would be as in Fig. 4b and 5 and also the evolution of the cell volume ratios would be as in Fig. 4(b). However, the ECM turnover would be faster than that shown in Fig. 5.

5. Invasion of tissue

We now assume that the abnormal tissue is generated (inhomogeneously) only in a certain region of the domain. In order to focus on the effect of the sensing of compression on cell proliferation and ECM

production and degradation, we assume that the behaviour of normal and abnormal cells differs only for the coefficients involving the sensing of stress. Therefore, $\gamma_n = \gamma_a = \gamma$ and $\delta_n = \delta_a = \delta$. The non-dimensional model is then given by

$$\begin{cases} \frac{\partial n}{\partial \tilde{t}} = \nabla \cdot (n \tilde{\Sigma}'(\psi) \nabla \psi) + (H_\sigma(\psi - \psi_n) - \tilde{\delta})n, \\ \frac{\partial a}{\partial \tilde{t}} = \nabla \cdot (a \tilde{\Sigma}'(\psi) \nabla \psi) + (H_\sigma(\psi - \psi_a) - \tilde{\delta})a, \\ \frac{\partial m_n}{\partial \tilde{t}} = \tilde{\mu}_n n - \tilde{v} \tilde{c} m_n, \\ \frac{\partial m_a}{\partial \tilde{t}} = \tilde{\mu}_a a - \tilde{v} \tilde{c} m_a, \\ \frac{\partial \tilde{c}}{\partial \tilde{t}} = \tilde{\kappa} \nabla^2 \tilde{c} + \frac{1}{\tilde{\tau}}(n + \tilde{\pi}_a a - \tilde{c}), \end{cases} \quad (46)$$

where

$$\tilde{t} = \gamma t, \quad \tilde{\mathbf{x}} = \sqrt{\frac{\gamma}{KE}} \mathbf{x}, \quad \tilde{c} = \frac{c}{\pi_n \tau} \quad (47)$$

and where

$$\begin{aligned} \tilde{\mu}_i &= \frac{\mu_i}{\gamma}, \quad i = n, a, \\ \tilde{\pi}_a &= \frac{\pi_a}{\pi_n}, \quad \tilde{\tau} = \gamma \tau, \quad \tilde{\kappa} = \frac{\kappa}{KE}, \quad \tilde{\delta} = \frac{\delta}{\gamma}, \quad \tilde{v} = \frac{v \pi_n \tau}{\gamma}, \quad \tilde{\Sigma}'(\psi) = \frac{\Sigma'(\psi)}{E}. \end{aligned} \quad (48)$$

We recall that n, a, m_n and m_a are already dimensionless and observe that $\tilde{v}/\tilde{\mu}_n = M_n$ and $\tilde{v}/\tilde{\mu}_a = M_a$ defined in (20) and (45), respectively. We observe that measurements of K and E are not available in the literature, to our knowledge. However, from the fact that the order of magnitude of the cell velocity is microns per hours and assuming a change of volume ratio of 0.1 over the tumour diameter, we can argue that the order of magnitude of the product KE is $1 \mu\text{m}^2/\text{s}$. Hence, the characteristic length in (46) is of the order of $100 \mu\text{m}$ and $\tilde{\kappa}$ is of the order of 100.

Mathematically, the domain of integration in the following simulations is the entire space. However, when for numerical reasons we restrict to a finite domain, zero-flux boundary conditions for the cells and the MDEs are imposed on the boundaries of the domain.

To start with, we consider the spatially dependent version of the problem dealt with at the beginning of Section 4.3. As we are working with the same assumption implying (31), i.e. the sum of m_n and m_a is constant, the evolution equations for the cells decouple from those for the ECM, and we can focus on

$$\begin{cases} \frac{\partial n}{\partial \tilde{t}} = \nabla \cdot [n \tilde{\Sigma}'(n+a) \nabla(n+a)] + [H_\sigma(n+a - \psi_n) - \tilde{\delta}]n, \\ \frac{\partial a}{\partial \tilde{t}} = \nabla \cdot [a \tilde{\Sigma}'(n+a) \nabla(n+a)] + [H_\sigma(n+a - \psi_a) - \tilde{\delta}]a, \end{cases} \quad (49)$$

with no-flux boundary conditions and the initial conditions

$$a(\tilde{t} = 0, \tilde{x}) = a_0 \exp\{-\tilde{x}^2/2\sigma_0\} \quad \text{and} \quad n(\tilde{t} = 0, \tilde{x}) = \psi_n - M_n - a_0 \exp\{-\tilde{x}^2/2\sigma_0\},$$

such that the overall volume ratio corresponds to the homogeneous stationary configuration.

Initially, the same process occurring in the relaxed replacement phase (Phase I) occurs, with the abnormal population partially replacing the normal one (see Fig. 7), i.e.

$$a(\tilde{t}, \tilde{x}) = a_0 e^{-\tilde{x}^2/2\sigma_0} e^{(1-\tilde{\delta})\tilde{t}} \quad \text{and} \quad n(\tilde{t}, \tilde{x}) = \psi_n - M_n - a(\tilde{t}, \tilde{x}), \quad (50)$$

until

$$\tilde{t} \approx \frac{1}{1-\tilde{\delta}} \log \frac{\tilde{\delta}(\psi_n - M_n)}{a_0}, \quad (51)$$

for small σ . The simulation presented in Fig. 7 (which refers to (49)) follows the estimate above. In fact, for the values of the parameters used, (51) predicts that compression starts for $\tilde{t} > 4.1$ and coherently in Fig. 7(c) the line corresponding to $\tilde{t} = 5$ barely departs from the initial value. At this point in

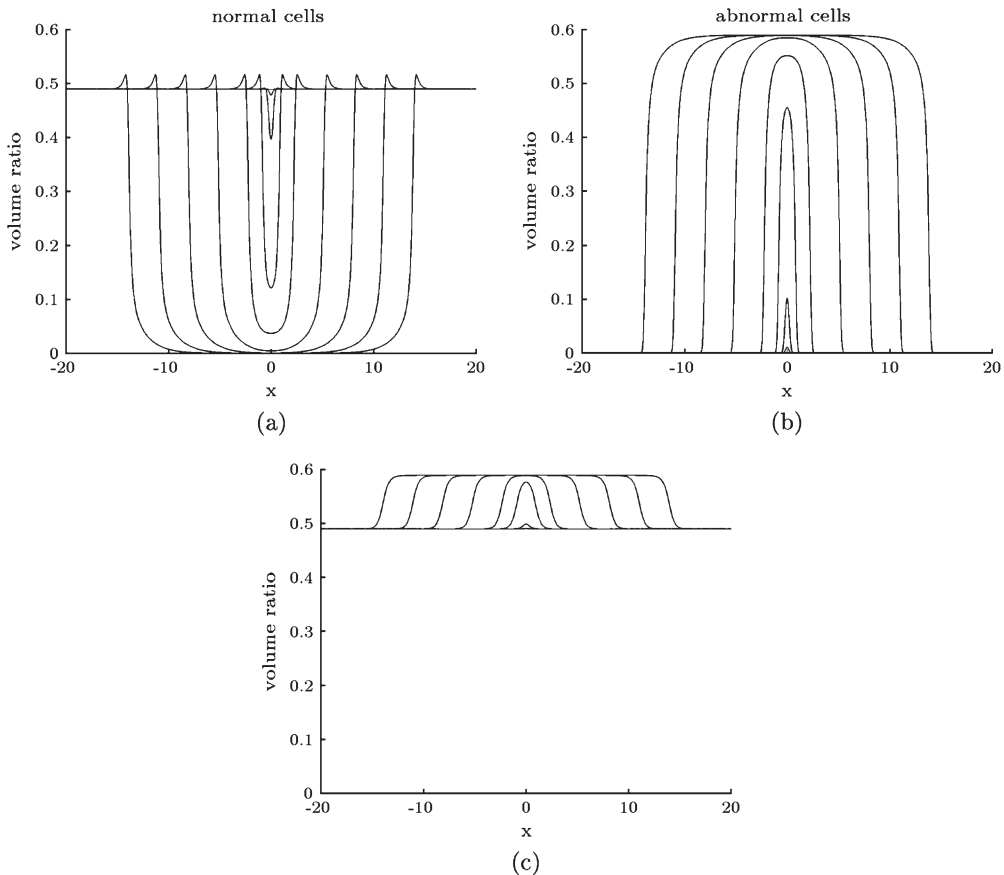


FIG. 7. Tissue invasion for the following set of parameters: $\tilde{\delta} = 0.1$, $\psi_n = 0.6$, $\psi_a = 0.8$, $\sigma = 0.1$, $M_n = 0.2$, and at times $\tilde{t} = 0, 2.5, 5, 10, 20, 40, 60, 80$. At smaller times, the dynamics follows the one described in Section 4.3. At later times, the travelling wave characteristic and the transition layer are evident. The compression of the normal tissue due to the expansion of the hyperplasia is also put in evidence by the peaks in (a) (see cartoon in Fig. 4(a)). In (c), the overall cell volume ratio ($n+a$) at early times keeps its initial value following the relaxed replacement phase depicted in Fig. 4(a). We recall that the unit dimensionless distance corresponds to a dimensional distance of the order of 0.1 mm.

$\tilde{x} = 0$, compression starts and cells start moving away from the compressed regions. So, Phase II and Phase III are different than in the spatially independent case treated in the previous sections. In the overcompressed regions where $\psi > \psi_n$, something like Phase II occurs with n exponentially decaying to zero and replaced by a . In the simulation, the maximum overall compression characterizing Phase III is achieved after about 40 (days, if $\gamma = 1$ per days), corresponding to a size of the hyperplasia which depends on \sqrt{KE} and can be evaluated of the order of 1 mm. This means that if the product of the permeability coefficient and the ‘Young’s modulus’ increases by two order of magnitude, i.e. for more permeable tissues and stiffly reacting cells, then the size needed to reach the overcompressed status increases by an order of magnitude. Of course, if growth is too slow with respect to motion, then the time needed to reach this phase increases further.

It is interesting to note that, as observed experimentally, the growth of the hyperplasia is accompanied by a compression of the normal cells near the interface separating the two tissues as shown by the peaks in Fig. 7(a) and qualitatively by the cartoon in Fig. 4(a).

Recalling the characteristics of $H_\sigma(\phi)$ in (5) and in particular that it is non-constant for $\phi \in [0, \sigma]$, in the hyperplastic tissue the solution tends to achieve an overall volume ratio in the interval $[\psi_a, \psi_a + \phi]$, while in the normal tissue it tends to keep an overall volume ratio in the interval $[\psi_n, \psi_n + \phi]$. Actually, from the maximum principle, $n + a + M_n \in [\psi_n, \psi_a + \phi]$ (see Fig. 7(c)).

This is true even if one starts from a segregated initial condition, i.e.

$$\begin{aligned} n(t = 0, \mathbf{x}) = 0 \quad \text{and} \quad a(t = 0, \mathbf{x}) \neq 0 \quad \text{in } \Omega, \\ n(t = 0, \mathbf{x}) \neq 0 \quad \text{and} \quad a(t = 0, \mathbf{x}) = 0 \quad \text{in } \mathbb{R}^3 - \Omega. \end{aligned} \quad (52)$$

In this case, because of the degenerate character of the parabolic equations, the solution for a will always have compact support with a sharp front $\mathbf{x}_T(t)$ separating the two tissues. Coherently, the two populations move with the same velocity at the interface

$$\tilde{\mathbf{v}} = -\mathbf{n} \cdot \nabla n = -\mathbf{n} \cdot \nabla a \quad \text{on } \partial\Omega. \quad (53)$$

On the other hand, though the equation governing the evolution of $n + a$ is still degenerate parabolic, its solution is regular everywhere because the solution stays away from zero. In fact, in addition to (53), one has

$$n = a \quad \text{on } \partial\Omega. \quad (54)$$

This leads to the formation of a transition layer near the interface linking the solution with $n + a = a \in [\psi_a, \psi_a + \sigma]$ in the abnormal tissue and that with $n + a = n \in [\psi_n, \psi_n + \sigma]$ in the normal tissue.

The evolution shown in Fig. 7 suggests the existence of a travelling wave solution. In order to identify the velocity of invasion of the hyperplasia and to describe the structure of the transition layer, we will look for a solution in the form of a travelling wave in the limit of small σ . Assuming $\tilde{\psi}_a - \tilde{\psi}_n \ll \tilde{\psi}_n$ where, for instance, $\tilde{\psi}_n = \psi_n - M_n$, and introducing the perturbations $\hat{n} = n - \tilde{\psi}_n$ and $\hat{a} = a - \tilde{\psi}_n$ one can write to leading order the following problem:

$$\begin{cases} -\tilde{v}\hat{a}' = \tilde{\psi}_n\hat{a}'' + (1 - \tilde{\delta})\tilde{\psi}_n, & z < 0, \\ -\tilde{v}\hat{n}' = \tilde{\psi}_n\hat{n}'' - \tilde{\delta}\tilde{\psi}_n, & z > 0, \\ \hat{a} = \psi_a - \psi_n & \text{at } z = z_1, \\ \hat{n} = 0 & \text{at } z = z_2, \\ \hat{n} = \hat{a} \quad \text{and} \quad \hat{n}' = \hat{a}' = -\tilde{v} & \text{at } z = 0, \end{cases} \quad (55)$$

where the primes denote derivation with respect to $z = \tilde{x} - \tilde{v}\tilde{t}$, \tilde{v} is the velocity of the propagating front and $z \in [z_1, z_2]$ denotes the transition layer.

In order to identify z_1, z_2 we also assume

$$\begin{aligned}\hat{a}' &= 0 & \text{at } z &= z_1, \\ \hat{n}' &= 0 & \text{at } z &= z_2.\end{aligned}\tag{56}$$

One then has

$$\begin{aligned}\hat{a} &= \tilde{\psi}_a - \tilde{\psi}_n + \frac{(\tilde{\delta} - 1)\tilde{\psi}_n}{\tilde{v}}(z - z_1) + \frac{(\tilde{\delta} - 1)\tilde{\psi}_n^2}{\tilde{v}^2} \left\{ \exp\left[-\frac{\tilde{v}(z - z_1)}{\tilde{\psi}_n}\right] - 1 \right\} & \text{for } z \in [z_1, 0], \\ \hat{n} &= \frac{\tilde{\delta}\tilde{\psi}_n}{\tilde{v}}(z - z_2) + \frac{\tilde{\delta}\tilde{\psi}_n^2}{\tilde{v}^2} \left\{ \exp\left[-\frac{\tilde{v}(z - z_2)}{\tilde{\psi}_n}\right] - 1 \right\} & \text{for } z \in [0, z_2],\end{aligned}\tag{57}$$

where

$$\begin{aligned}z_1 &= \frac{\tilde{\psi}_n}{\tilde{v}} \ln\left(1 - \frac{\tilde{v}^2}{\tilde{\psi}_n(1 - \tilde{\delta})}\right), \\ z_2 &= \frac{\tilde{\psi}_n}{\tilde{v}} \ln\left(1 + \frac{\tilde{v}^2}{\tilde{\psi}_n\tilde{\delta}}\right),\end{aligned}\tag{58}$$

and the velocity of propagation is implicitly given by

$$\tilde{\psi}_a - \tilde{\psi}_n = -\frac{\tilde{\psi}_n^2}{\tilde{v}^2} \left[\tilde{\delta} \ln\left(1 + \frac{\tilde{v}^2}{\tilde{\psi}_n\tilde{\delta}}\right) + (1 - \tilde{\delta}) \ln\left(1 - \frac{\tilde{v}^2}{\tilde{\psi}_n(1 - \tilde{\delta})}\right) \right].\tag{59}$$

In the limit $\tilde{v} \ll 1$, corresponding to $\psi_a \approx \psi_n$, one has the following approximations for the transition layer width:

$$z_2 - z_1 \approx \sqrt{\frac{2(\psi_a - \psi_n)}{\tilde{\delta}(1 - \tilde{\delta})}},\tag{60}$$

and for the velocity of propagation

$$\tilde{v} \approx \sqrt{2\tilde{\delta}(1 - \tilde{\delta})(\psi_a - \psi_n)},\tag{61}$$

which in dimensional form writes

$$v \approx \sqrt{2KE\delta\left(1 - \frac{\delta}{\gamma}\right)(\psi_a - \psi_n)}.\tag{62}$$

Figure 8 shows a comparison between the travelling wave solution (57) and the one obtained numerically. Negative values of z correspond to the abnormal tissue and positive values to the normal tissue. Hence, $z = 0$ is the interface between the tissues. Figure 9 compares the theoretical values of \tilde{v} , z_1 , z_2 , and then of the transition layer thickness with the one obtained from the simulations.

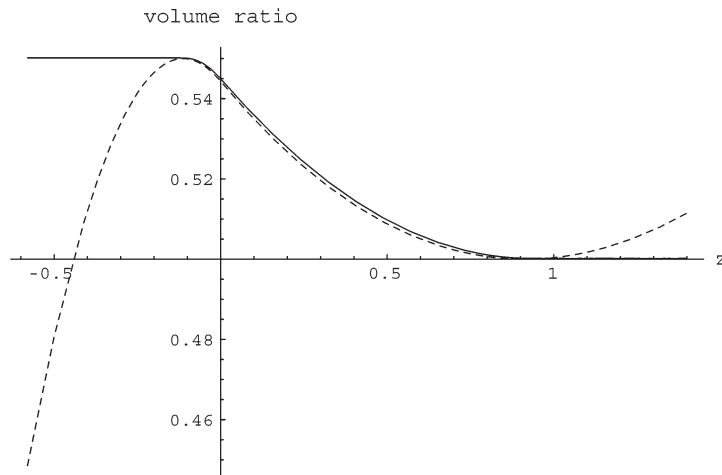


FIG. 8. Travelling wave solution. Comparison between numerical (full line) and analytical results (dashed line) given by (57) for $\tilde{\delta} = 0.1$, $\tilde{\psi}_n = 0.5$ and $\tilde{\psi}_a = 0.55$. $z = 0$ represents the moving interface between the normal and the abnormal tissue (on its right and left, respectively).

The simulation presented in Fig. 10 instead refers to the numerical solution of the full system of partial differential equations (46). In this case, the normal ECM is constantly digested and produced by the cells, but as the hyperplastic tissue is replacing the normal one, the abnormal ECM is now also replacing the normal one, still keeping the total amount of ECM constant, as explained in the space independent case discussed in Section 3. Also in this case, the travelling wave characteristic of the solution is evident. The evolution of the normal and abnormal cell populations are as in Fig. 7.

The concentration of MDEs increases because of the increase in the density of cells. Dynamics such as those discussed in Section 4.3 would be obtained for different values of the production rates of ECM by the two populations of cells, with the formation of an abnormal matrix which fills more space than normal.

6. Discussion and conclusions

In this paper, we have presented and developed a mathematical model focussing on the potential role that an incorrect mechano-transduction may play in causing the hyper-proliferation of cells in a tissue. In particular, we have shown how even a small loss of contact inhibition to growth and of compression responsiveness can represent by itself a clonal advantage which can give rise to the formation of hyperplasia and tumour growth. In replacing the normal tissue with abnormal tissue, we have identified three phases which qualitatively correspond to the phenomenological description given by biologists: a first phase in which some abnormal cells simply replace the normal ones, a second phase in which the hyper-proliferation of the abnormal cells causes a progressive compression within the tissue itself and a third phase in which the tissue reaches a compressed state, which presses on the surrounding environment. The difference in the ability to withstand the compressed state leads to a progressive invasion of the surrounding tissue.

At the same time, because of the constant remodelling of the ECM, the components produced by the tumour cells progressively replace those previously present. In the case in which the production

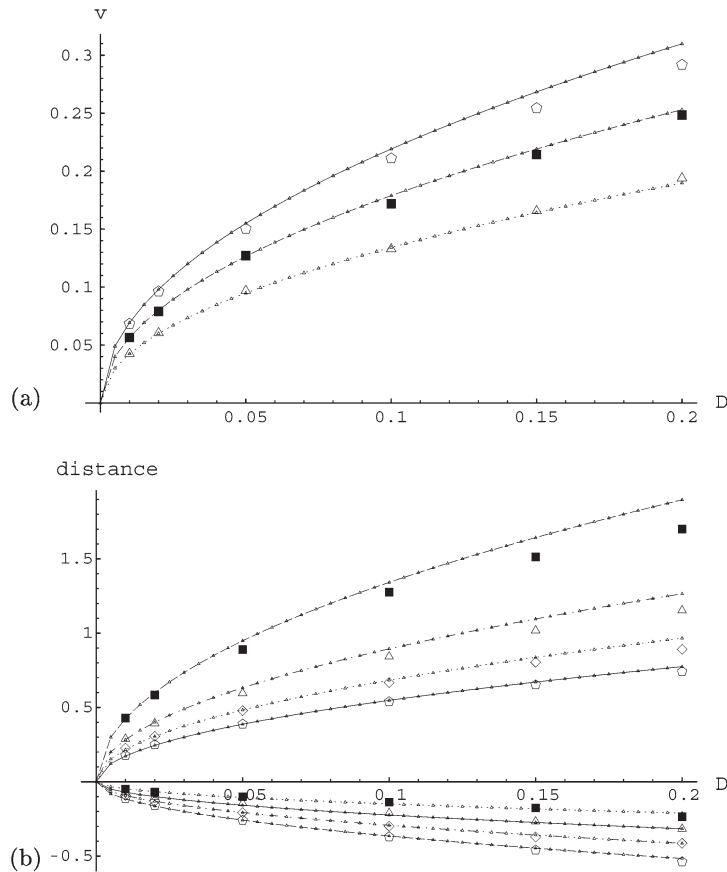


FIG. 9. Travelling wave solution. Comparison between analytical estimates (lines) and numerical results (symbols). (a) Velocity of propagation as a function of $D = \psi_a - \psi_n$ for $\tilde{\delta} = 0.1, 0.2, 0.4$ (from lower to upper curve). (b) Transition layer thickness. The lower family of curves refer to z_1 and the upper family to z_2 . The four curves refer to $\tilde{\delta} = 0.1, 0.2, 0.3, 0.4$ from above to below in both cases. The thickness of the transition layer is then given by the distance between the two curves for given values of the parameters.

rates or the release of MDEs is different than those corresponding to the physiological state, the model foresees the formation of a tissue characterized by a pathological percentage of ECM, as observed in many tumours and illnesses (see Fig. 6). For instance, an increased production of ECM not balanced by an increased release of MDEs leads to tissues with more ECM than normal, as in many fibrosis and in colon and prostate cancers. On the other hand, excessive degradation of ECM due to excessive production of MDEs leads to tissues with less ECM than normal, as in other malignant tumours and many other chronic inflammatory diseases.

From a travelling wave analysis exploiting the working assumption of only a slight incorrect sensing of the compression level, we obtained an estimate of the velocity of growth of the hyperplasia which is proportional to $\sqrt{\psi_a - \psi_n}$, i.e. to the square root of the difference of the volume ratios at which abnormal and normal cells stop duplicating by contact inhibition. Therefore, as expected, smaller disfunctions lead to slower invasions of the surrounding healthy tissue.

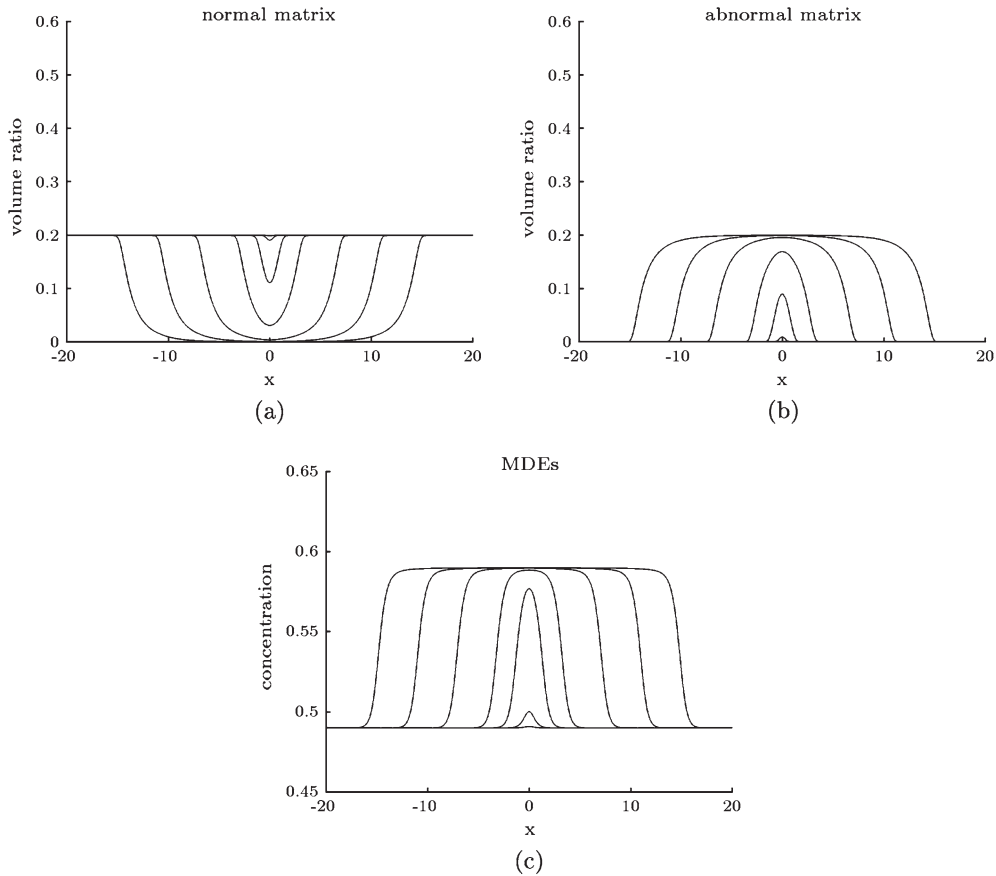


FIG. 10. Tissue invasion with degradation and renewal of ECM through the constant production of MDEs. The evolution of the cell populations are similar to those in Fig. 7. The values of the parameters are the same as in Fig. 7 and $\tilde{\pi}_a = 1$, $\tilde{\mu}_n = \tilde{\mu}_a = 0.1$, $\tilde{\tau} = 0.005$, $\tilde{\kappa} = 100$, $\tilde{\nu} = 0.5$. The cleavage of normal ECM inside the hyperplasia and its substitution with the one produced by abnormal cells follows the front of the hyperplasia. We recall that the unit dimensionless distance corresponds to a dimensional distance of the order of 0.1 mm.

Throughout the paper the role of nutrients and growth factors was neglected, though we are well aware of their importance, and that even the switching on and off of the duplication process is mediated by growth promoting and inhibitory factors, through the activation of suitable protein cascades.

In this respect, it would be interesting to develop a multiscale model which fully takes into account the cascade of events recalled in Fig. 2 possibly joined with those involving growth factors. Of course, in order to do that one needs to have estimates of the affinity constants. In absence of such measurements, one could initially start with a Boolean reasoning, which leads, for instance, to the assumption that a chronic cadherin under-expression leads to uncontrolled growth at any compression level (i.e. $\psi_a = 1$ in (15)).

Another interesting generalization of the model would be to take into account the fact that the remodelling process is strongly affected by the stresses and strains the tissue is subject to. Of course, in order to do that, quantitative measurement of the dependence of the production of ECM and MDEs

from stress and/or strain would be needed. Quantitative measurements on the growth rates and on the production rates of ECM and MDEs under different conditions would also allow a further validation of the model presented here.

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